# Vibration analysis of angle-ply laminates composite plate under thermomechanical effect 

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#### Abstract

ABSRACT

The paper presents mainly the dynamic response of an angle ply composite laminated plates subjected to thermo-mechanical loading. The response are analyzed by analytically using Newmark direct integration method with Navier solution, numerically by ANSYS. The experimental investigation is to fabricate the laminates and to find mechanical and thermal properties of glass-polyester such as longitudinal, transverse young modulus, shear modulus, longitudinal and transverse thermal expansion. Present of temperature could increase dynamic response of plate also depending on lamination angle, type of mechanical load and the value of temperature.


Keywords: composite laminated plate, dynamic response, thermo-mechanical loading


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الخلاصة هذه الدراسة ركزت بشكل رئيسي على الاستجابة الديناميكية للصفائح المركبة ذات زاوية معينة معرضة الى تاثير
درجة الحرارة. ويتم تحليل الاستجابة الديناميكية عن طريق التحليل النظري باستخدام طريقة نيومارك مع حل نافيير ، النحليل
العددي باستخدام برنامج .ANSYS في الجانب العملي تم تصنيع الصفائح المركبة المصنو عة من الالياف الزجاجية والبوليستر
لايجاد الخواص الحرارية والميكانيكية مثل معامل يونك الطولي ولئلي والعرضي و معامل القص والتمدد الحراري الطولي
والعرضي ـ وكان الاستتناج الرئبسي ان وجود درجة الحرارة بوكن ان يزيد الاستجابة الديناميكية اعتمادا على زاوية التصفيح
ونو الحرارة. الحيكانيكي
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## 1. INTRODUCTION

During the last decades, needs for composite materials consist of two or more types of materials mixed together homogenously have appeared to produce desirable properties, the constituents are combined at a macroscopic level, one constituent is called the reinforcing phase called fiber and the one in which is embedded is called the matrix. Reddy J.N.]
The analysis of structural vibration is necessary in order to calculate the natural frequencies of a structure, and the response to the expected excitation. In this way it can be determined whether a particular structure will fulfil its intended function and, in addition, the results of the dynamic loadings acting on a structure can be predicted, such as the dynamic stresses, fatigue life and noise levels. Ref. [Beards C.E]
Many researches had studied free vibration analysis and vibration of plate under mechanical or thermal or thermo-mechanical loading.
Chorng-Fuh Liu and Chih-Hsing Huang, 1996 performed a vibration analysis of laminated composite plates subjected to temperature change. The first order shear deformation theory of a plate is employed. The resulting finite element formulation leads to general nonlinear and coupled simulation equations and calculate the frequencies of vibration of a symmetric cross-ply plate. Hui-Shen Shen, et.al, 2003, studied the dynamic response of laminated plates subjected to thermo-mechanical loading and resting on a two-parameter elastic foundation. The formulation is based on higher order shear deformable plate theory and includes the thermal effect. Effects of foundation stiffness, thickness ratio, and temperature change on the dynamic response are discussed. Kullasup P. et al., 2010, analysed free vibration of symmetrically laminated composite plates with various boundary conditions by Kantorovich method. The beam function is used as an initial trial function in the repeated calculation, which is employed to calculate the natural frequency. Suresh K. J. et al., 2011, developed an analytical procedure is to evaluate the free vibration characteristics of laminated composite plates based on higher order shear deformation with zig-zag function. Slope discontinuities improved by Zig-zag function at the interfaces of laminated composite plates. The solutions are obtained using Navier's method.
Junaid Kameran Ahmed et al., 2013, presented a static and dynamic analysis of Graphite /Epoxy composite plates. In this work the behavior of laminated composite plates under transverse loading using an eight-node diso-parametric quadratic element based on First Order Shear Deformation Theory was studied,. Pushpendra k. kushwaha1 and jyoti vimal, 2014, the natural frequencies and mode shapes are compared for different boundary condition. Comparisons are made with the result for thin and thick composite laminated plate. Numerical results have been computed for the effect of number of layers, thickness ratio of plate, different boundary conditions, different aspect ratio, and different angle of fiber orientation of laminated composite plate.
The point of originality of the present work is how to derive the analytical solution of dynamic response for composite laminated plates by classical laminated plate theory for the first time under thermo-mechanical loading, by applied different type of loading on the symmetric and anti- symmetric angle ply composite laminated plates using Newmark direct integration method with Navier solution. Thermal and mechanical properties for composite plate made from (glasspolyester) with fiber volume fraction (0.3) are determined experimentally.
Also Finite element coded by ANSYS14.0 used to find natural frequency of composite laminate plate.

## 2. ANALYTICAL SOLUTION (CLASSICAL LAMINATE PLATE THEORY)

### 2.1 Displacement

Classical lamination theory (CLPT) based on the Kirchhoff hypothesis based on assuming the straight line perpendicular to the mid surface before deformation remains straight after deformation which means neglecting shear strains and transverse normal strain and stress in the analysis of laminated composite plates. Ref.[ Reddy J.N.]
$u(x, y, t)=u_{o}(x, y, t)-z \frac{\partial w_{0}}{\partial x}$
$v(x, y, t)=v_{o}(x, y, t)-z \frac{\partial w_{0}}{\partial y}$
$w(x, y, t)=w_{o}(x, y)$

Where $\frac{\partial w_{0}}{\partial x}, \frac{\partial w_{0}}{\partial y}$ denote the rotations about y and x axis respectively.
$u_{o}, v_{o}$ and $w_{o}$ denote the displacement components along ( $\mathrm{x}, \mathrm{y}, \mathrm{z}$ ) directions respectively of a point on the mid-plane (i.e.... $\mathrm{z}=0$ ).

### 2.2 Stress and Strain

The total strains can be written as follows

$$
\left\{\begin{array}{l}
\varepsilon_{\mathrm{xx}}  \tag{2.a}\\
\varepsilon_{\mathrm{yy}} \\
\gamma_{\mathrm{xy}}
\end{array}\right\}=\left\{\begin{array}{l}
\varepsilon_{\mathrm{xx}}^{(0)} \\
\varepsilon_{\mathrm{yy}}^{(0)} \\
\gamma_{\mathrm{xy}}^{(0)}
\end{array}\right\}+\mathrm{z} *\left\{\begin{array}{c}
\varepsilon_{\mathrm{xx}}^{(1)} \\
\varepsilon_{\mathrm{yy}}^{(1)} \\
\gamma_{\mathrm{xy}}^{(1)}
\end{array}\right\}=\left\{\begin{array}{c}
\frac{\partial u_{0}}{\partial x}-\alpha_{x x} \Delta T \\
\frac{\partial v_{0}}{\partial y}-\alpha_{y y} \Delta T \\
\frac{\partial u_{0}}{\partial y}+\frac{\partial v_{0}}{\partial x}-\alpha_{x y} \Delta T
\end{array}\right\}+z\left\{\begin{array}{c}
-\frac{\partial^{2} \mathrm{w}_{\mathrm{o}}}{\partial \mathrm{x}^{2}} \\
-\frac{\partial^{2} \mathrm{w}_{\mathrm{o}}}{\partial \mathrm{y}^{2}} \\
-2 * \frac{\partial^{2} \mathrm{w}_{\mathrm{o}}}{\partial \mathrm{x} \partial \mathrm{y}}
\end{array}\right\}
$$

Where $\left(\varepsilon_{x x}^{(0)}, \varepsilon_{y y}^{(0)}, \gamma_{x y}^{(0)}\right)$ are the membrane strains and $\left(\varepsilon_{x x}^{(1)}, \varepsilon_{y y}^{(1)}, \gamma_{x y}^{(1)}\right)$ are the flexural (bending) strains, known as the curvatures $\alpha_{x x}, \alpha_{y y}$ and $\alpha_{x y}$ are thermal expansion coefficients defined

$$
\begin{equation*}
\alpha_{x x}=\alpha_{11}(\cos \theta)^{2}+\alpha_{22}(\sin \theta)^{2} \tag{2.b}
\end{equation*}
$$

$\alpha_{y y}=\alpha_{11}(\sin \theta)^{2}+\alpha_{22}(\cos \theta)^{2}$
$2 \alpha_{x y}=2\left(\alpha_{11}-\alpha_{22}\right) \sin \theta \cos \theta$
$\alpha_{11}$ and $\alpha_{22}$ Are longitudinal and transverse thermal expansions respectively. And $\theta$ is the lamination angle.

The change in temperature defined
$\Delta T=$ applied temperature - reference temperature

Where reference temperature $T_{r e f}=25 C^{\circ}$ [Reddy J.N.]. The transformed stress-strain relations of an orthotropic lamina in a plane state of stress are; for $\bar{Q}_{i j}$ see [Reddy J.N.]

$$
\left\{\begin{array}{c}
\sigma_{x x}  \tag{3}\\
\sigma_{y y} \\
\sigma_{x y}
\end{array}\right\}_{k}=\left[\begin{array}{lll}
\bar{Q}_{11} & \bar{Q}_{12} & \bar{Q}_{16} \\
\bar{Q}_{12} & \bar{Q}_{22} & \bar{Q}_{26} \\
\bar{Q}_{16} & \bar{Q}_{26} & \bar{Q}_{66}
\end{array}\right]_{k}\left\{\begin{array}{c}
\varepsilon_{x x-}-\alpha_{x x} \Delta T \\
\varepsilon_{y y-\alpha_{y y} \Delta T} \\
\gamma_{x y-2 \alpha_{x y} \Delta T}
\end{array}\right\}
$$

The resultant of inplane force $N x x$, Nyy and $N x y$ and moments Mxx, Myy and Mxy acting on a laminate can be obtained from integration of the stress in each layer or lamina through the laminate thickness. Knowing the stressed in terms of the displacements, the inplane force resultants Nxx, Nyy, Nxy, Mxx, Myy and Mxy can be obtained.

The inplane force resultants are defined as
$\left\{\begin{array}{l}\mathrm{N}_{\mathrm{xx}} \\ \mathrm{N}_{\mathrm{yy}} \\ \mathrm{N}_{\mathrm{xy}}\end{array}\right\}=\sum_{\mathrm{k}=1}^{\mathrm{N}} \int_{\mathrm{z}_{\mathrm{k}}}^{\mathrm{z}_{\mathrm{k}+1}}\left\{\begin{array}{l}\sigma_{\mathrm{xx}} \\ \sigma_{\mathrm{yy}} \\ \sigma_{\mathrm{xy}}\end{array}\right\}_{\mathrm{k}} \mathrm{dz}$

Where $\sigma_{x}, \sigma_{y}$ and $\sigma_{x y}$ are normal and shear stress.

$$
\left\{\begin{array}{l}
N_{x x}  \tag{4.b}\\
N_{y y} \\
N_{x y}
\end{array}\right\}=\left[\begin{array}{lll}
A_{11} & A_{12} & A_{16} \\
A_{12} & A_{22} & A_{26} \\
A_{16} & A_{26} & A_{66}
\end{array}\right]\left\{\begin{array}{c}
\varepsilon_{x y}^{0} \\
\varepsilon_{x y}^{0} \\
\gamma_{x y}^{0}
\end{array}\right\}+\left[\begin{array}{lll}
B_{11} & B_{12} & B_{16} \\
B_{12} & B_{22} & B_{26} \\
B_{16} & B_{26} & B_{66}
\end{array}\right]\left\{\begin{array}{c}
\varepsilon_{x x}^{1} \\
\varepsilon_{y y}^{1} \\
\gamma_{x y}^{1}
\end{array}\right\}-\left\{\begin{array}{l}
N_{x x}^{t} \\
N_{y y}^{t} \\
N_{x y}^{t}
\end{array}\right\}
$$

$\left\{\begin{array}{l}M_{x x} \\ M_{y y} \\ M_{x y}\end{array}\right\}=\sum_{k=1}^{N} \int_{z_{k}}^{z_{k+1}}\left\{\begin{array}{l}\sigma_{x x} \\ \sigma_{y y} \\ \sigma_{x y}\end{array}\right\}_{k} z d z$
$\left\{\begin{array}{l}M_{x x} \\ M_{y y} \\ M_{x y}\end{array}\right\}=\left[\begin{array}{lll}B_{11} & B_{12} & B_{16} \\ B_{12} & B_{22} & B_{26} \\ B_{16} & B_{26} & B_{66}\end{array}\right]\left\{\begin{array}{c}\varepsilon_{x y}^{0} \\ \varepsilon_{x y}^{0} \\ \gamma_{x y}^{0}\end{array}\right\}+\left[\begin{array}{lll}D_{11} & D_{12} & D_{16} \\ D_{12} & D_{22} & D_{26} \\ D_{16} & D_{26} & D_{66}\end{array}\right]\left\{\begin{array}{c}\varepsilon_{x x}^{1} \\ \varepsilon_{y y}^{1} \\ \gamma_{x y}^{1}\end{array}\right\}-\left\{\begin{array}{l}M_{x x}^{t} \\ M_{y y}^{t} \\ M_{x y}^{t}\end{array}\right\}$

Here, $A_{i j}$ are the extensional stiffness, $B_{i j}$ the coupling stiffness, and $D_{i j}$ the bending stiffness.
$A_{i j}=\sum_{k=1}^{N}\left(\bar{Q}_{i j}\right)_{k}\left(z_{k+1}-z_{k}\right)$
$B_{i j}=\frac{1}{2} \sum_{k=1}^{N}\left(\bar{Q}_{i j}\right)_{k}\left(z^{2}{ }_{k+1}-z^{2}{ }_{k}\right)$
$D_{i j}=\frac{1}{3} \sum_{k=1}^{N}\left(\bar{Q}_{i j}\right)_{k}\left(z^{3}{ }_{k+1}-z^{3}{ }_{k}\right)$

And Where $\left\{N^{t}\right\}$ and $\left\{M^{t}\right\}$ are thermal stress and bending results, respectively

$$
\left\{\begin{array}{l}
N_{x x}^{t}, M_{x x}^{t}  \tag{6.d}\\
N_{y y}^{t}, M_{y y}^{t} \\
N_{x y}^{t}, M_{x y}^{t}
\end{array}\right\}=\sum_{k=1}^{N} \int_{-h / 2}^{h / 2}\left[\begin{array}{ccc}
\bar{Q}_{11} & \bar{Q}_{12} & \bar{Q}_{16} \\
\bar{Q}_{12} & \bar{Q}_{22} & \bar{Q}_{26} \\
\bar{Q}_{16} & \bar{Q}_{26} & \bar{Q}_{66}
\end{array}\right]\left\{\begin{array}{c}
\alpha_{x x} \\
\alpha_{y y} \\
2 \alpha_{x y}
\end{array}\right\}(1, z) \Delta T d z
$$

### 2.3 Equation of Motion

The equations of motion are obtained by setting the coefficient of $\delta u_{0}, \delta v_{0}, \delta w_{0}$ to zero separately
$\frac{\partial N_{x x}}{\partial x}+\frac{\partial N_{x y}}{\partial y}=I_{0} \frac{\partial^{2} u}{\partial t^{2}}-I_{1} \frac{\partial^{3} w}{\partial x \partial t^{2}}$
$\frac{\partial N_{x y}}{\partial x}+\frac{\partial N_{y y}}{\partial y}=I_{0} \frac{\partial^{2} v}{\partial t^{2}}-I_{1} \frac{\partial^{3} w}{\partial y \partial t^{2}}$
$\frac{\partial^{2} M_{x x}}{\partial x^{2}}+2 \frac{\partial^{2} M_{x y}}{\partial x \partial y}+\frac{\partial^{2} M_{y y}}{\partial y^{2}}+\widehat{N}_{x x} \frac{\partial^{2} w}{\partial x^{2}}+\widehat{N}_{y y} \frac{\partial^{2} w}{\partial y^{2}}+\widehat{N}_{x y} \frac{\partial^{2} w}{\partial x \partial y}=I_{0} \frac{\partial^{2} w}{\partial t^{2}}+I_{1}\left(\frac{\partial^{3} u}{\partial x \partial t^{2}}+\frac{\partial^{3} v}{\partial y \partial t^{2}}\right)-$
$I_{2}\left(\frac{\partial^{4} w}{\partial x^{2} \partial t^{2}}+\frac{\partial^{4} w}{\partial y^{2} \partial t^{2}}\right)-q(x, y, t)$

Where

$$
\begin{equation*}
\left(I_{0}, I_{1} \cdot I_{2}\right)=\sum_{k=1}^{N} \int_{z_{k-1}}^{z_{k}} \rho^{(k)}\left(1, z, z^{2}\right) d z \tag{8}
\end{equation*}
$$

$\rho^{(\mathrm{k})}$ being the material density of $\mathrm{k}^{\text {th }}$ layer and $\mathrm{q}(\mathrm{x}, \mathrm{y}, \mathrm{t})$ is a dynamic force subjected on a system. $\widehat{N}_{x x}, \widehat{N}_{y y}$ and $\widehat{N}_{x y}$ equal to zero because there were no buckling.

These equations of motion ( $7 \mathrm{a}-\mathrm{c}$ ) can be expressed in terms of displacements ( $\delta u_{0}, \delta v_{0}, \delta w_{0}$ ) by substituting the forces results from Eqs. $(4,5,8)$ into Eq. (7.a) to(7.c) and get partial differential equations,

$$
\begin{align*}
& {\left[\begin{array}{lll}
c_{11} & c_{12} & c_{13} \\
c_{12} & c_{22} & c_{23} \\
c_{13} & c_{23} & c_{33}
\end{array}\right]\left\{\begin{array}{l}
u_{0} \\
v_{0} \\
w_{0}
\end{array}\right\}+\left[\begin{array}{ccc}
m_{11} & 0 & 0 \\
0 & m_{22} & 0 \\
0 & 0 & m_{33}
\end{array}\right]\left\{\begin{array}{c}
\ddot{u}_{0} \\
\ddot{v}_{0} \\
\ddot{w}_{0}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0 \\
q
\end{array}\right\}+\left\{\begin{array}{l}
f_{1}^{t} \\
f_{2}^{t} \\
f_{3}^{t}
\end{array}\right\}}  \tag{9.a}\\
& c_{11}=A_{11} d_{x}^{2}+2 A_{16} d_{x} d_{y}+A_{66} d_{y}^{2}  \tag{9.b}\\
& c_{12}=A_{16} d_{x}^{2}+\left(A_{12}+A_{66}\right) d_{x} d_{y}+A_{26} d_{y}^{2}  \tag{9.c}\\
& c_{13}=-\left[B_{11} d_{x}^{3}+3 B_{16} d_{x}^{2} d_{y}+\left(B_{12}+2 B_{66}\right) d_{x} d_{y}^{2}+B_{26} d_{y}^{3}\right] \tag{9.d}
\end{align*}
$$

$c_{22}=A_{66} d_{x}^{2}+2 A_{26} d_{x} d_{y}+A_{22} d_{y}^{2}$
$c_{23}=-\left[B_{16} d_{x}^{3}+\left(B_{12}+2 B_{66}\right) d_{x}^{2} d_{y}+3 B_{26} d_{x} d_{y}^{2}+B_{22} d_{y}^{3}\right]$
$c_{33}=-D_{11} d_{x}^{4}-4 D_{16} d_{x}^{3} d_{y}-2\left(D_{12}+2 D_{66}\right) d_{x}^{2} d_{y}^{2}-4 D_{26} d_{x}-D_{22} d_{y}^{4}-\left(A_{11} \alpha_{x x}+A_{12} \alpha_{y y}+\right.$ $\left.2 A_{16} \alpha_{x y}\right) \Delta T d_{x}^{2}-\left(A_{16} \alpha_{x x}+A_{26} \alpha_{y y}+4 A_{66} \alpha_{x y}\right) \Delta T d_{x} d_{y}-\left(A_{12} \alpha_{x x}+A_{22} \alpha_{y y}+\right.$ $\left.2 A_{26} \alpha_{x y}\right) \Delta T d_{y}^{2}$
$f_{1}^{t}=\frac{\partial N_{x x}^{t}}{\partial x}+\frac{\partial N_{x y}^{t}}{\partial y}$
$f_{2}^{t}=\frac{\partial N_{x y}^{t}}{\partial x}+\frac{\partial N_{y y}^{t}}{\partial y}$
$f_{3}^{t}=-\left(\frac{\partial^{2} M_{x x}^{t}}{\partial x^{2}}+2 \frac{\partial^{2} M_{x y}^{t}}{\partial y \partial x}+\frac{\partial^{2} M_{y y}^{t}}{\partial y^{2}}\right)$

And the coefficients $m_{i j}$ is defined by
$m_{11}=-I_{0} d_{t}^{2}, m_{13}=I_{1} d_{x} d_{t}^{2}, m_{22}=-I_{0} d_{t}^{2} ; m_{23}=I_{1} d_{y} d_{t}^{2}, m_{33}=I_{0} d_{t}^{2}-I_{2} d_{t}^{2}\left(d_{x}^{2}+d_{y}^{2}\right)$

To solve equation (9-a) used Navier solution with state space approach.
For angle-ply rectangular laminates with edges $\mathrm{y}=0$ and $\mathrm{y}=\mathrm{b}$ simply supported and the other two edges $x=0$ and $x=a$ simply supported. Assume the following representation of the displacement [Reddy .J.N.]
$u_{0}(x, y, t)=\sum_{n=1}^{\infty} \sum_{m=1}^{\infty} U_{m n}(t) \sin \alpha x \cos \beta y$
$v_{0}(x, y, t)=\sum_{n=1}^{\infty} \sum_{m=1}^{\infty} V_{m n}(t) \cos \alpha x \sin \beta y$
$w_{0}(x, y, t)=\sum_{n=1}^{\infty} \sum_{m=1}^{\infty} W_{m n}(t) \sin \alpha x \sin \beta y$
Where: $\quad \alpha=\frac{m \pi}{a} ; \quad \beta=\frac{n \pi}{a}$
$m=$ No. of the mode in x-direction $(m=1,2,3)$
$n=$ No. of the mode in $y$-direction $(n=1,2,3)$
$U_{m n}, V_{m n}, W_{m n}$ are coefficients to be determined; and
$\Delta T(x, y, t)=\sum_{n=1}^{\infty} \sum_{m=1}^{\infty} T_{m n}(t) \sin \alpha x \sin \beta y$

$$
\begin{equation*}
T_{m n}(t)=\frac{4}{a b} \int_{0}^{a} \int_{0}^{b} \Delta T(x, y, t) \sin \alpha x \sin \beta y d x d y \mathrm{By} \tag{11.b}
\end{equation*}
$$

substituting Eqs. (10 and 11) in partial differential Eq. (9.a) and the result
$\left[\begin{array}{lll}\hat{c}_{11} & \hat{c}_{12} & \hat{c}_{13} \\ \hat{c}_{12} & \hat{c}_{22} & \hat{c}_{23} \\ \hat{c}_{13} & \hat{c}_{23} & \hat{c}_{33}\end{array}\right]\left\{\begin{array}{l}u_{m n} \\ v_{m n} \\ w_{m n}\end{array}\right\}+\left[\begin{array}{ccc}\hat{m}_{11} & 0 & 0 \\ 0 & \widehat{m}_{22} & 0 \\ 0 & 0 & \widehat{m}_{33}\end{array}\right]\left\{\begin{array}{l}\ddot{u}_{m n} \\ \ddot{v}_{m n} \\ \ddot{w}_{m n}\end{array}\right\}=\left\{\begin{array}{c}0 \\ 0 \\ Q_{m n}\end{array}\right\}+\left\{\begin{array}{c}\alpha N_{m n}^{1} \\ \beta N_{m n}^{2} \\ -2 \alpha \beta M_{m n}^{6}\end{array}\right\}$
(12.a)

Where
$\hat{c}_{11}=\left(A_{11} \alpha^{2}+A_{66} \beta^{2}\right)$
(12.b)
$\hat{c}_{12}=\left(A_{12}+A_{66}\right) \alpha \beta$
(12.c)

$$
\begin{align*}
& \hat{c}_{13}=-\left(3 B_{16} \alpha^{2}+B_{26} \beta^{2}\right) \beta  \tag{12.d}\\
& \hat{c}_{22}=\left(A_{66} \alpha^{2}+A_{22} \beta^{2}\right)  \tag{12.e}\\
& \hat{c}_{23}=-\left(B_{16} \alpha^{2}+3 B_{26} \beta^{2}\right) \alpha  \tag{12.f}\\
& \hat{c}_{33}=D_{11} \alpha^{4}+2\left(D_{12}+2 D_{66}\right) \alpha^{2} \beta^{2}+D_{22} \beta^{4} \tag{12.g}
\end{align*}
$$

$\widehat{m}_{11}=\widehat{m}_{22}=I_{0}$
(12.h)
$\widehat{m}_{33}=\left(I_{0}+I_{2}\left(\alpha^{2}+\beta^{2}\right)\right)$
The dynamic load subjected on the system,
$q(x, y, t)=\sum_{n=1}^{\infty} \sum_{m=1}^{\infty} Q_{m n}(x, y, t) \sin \alpha x \sin \beta y$
(13.a)
$Q_{m n}(x, y, t)=\frac{4}{a b} \int_{0}^{b} \int_{0}^{a} q(x, y) \cdot \sin \frac{m \pi x}{a} \cdot \sin \frac{n \pi y}{b} \cdot d x \cdot d y \cdot f(t)$
(13.b)
$Q_{m n}(x, y, t)=\bar{q}(x, y) \cdot f(t)$

Many types of $q(x, y)$ loading can be considered as [Muhannad L. S. Al-Waily,2004]:

1. Uniformly distributed load $q_{0}$ at area of plate (a*b). By substituting the load into Eq. (13.b), gives:

$$
\begin{equation*}
\bar{q}(x, y)=\frac{16}{m n \pi^{2}} q_{0} \tag{14.a}
\end{equation*}
$$

2. Point load $\mathrm{P}_{\mathrm{o}}$ at $\mathrm{x}=\mathrm{a}_{1}$ and $\mathrm{y}=\mathrm{b}_{1}$. By substituting the load into Eq.(13.b), gives:
$\bar{q}(x, y)=\frac{4}{a b} P_{0} \sin \frac{m \pi a_{1}}{a} \sin \frac{n \pi b_{1}}{b}$
(14.b)
3. Sinusoidal distributed loading, $q(x, y)=q_{0} \sin \frac{\pi x}{a} \sin \frac{\pi y}{b}$. By substituting the load into Eq.
(13.b), gives:
$\bar{q}(x, y)=q_{0}$
(14.c)
4. Uniformly distributed load $\mathrm{q}_{\mathrm{o}}$ at central area ( $\mathrm{A} * \mathrm{~B}$ ). By substituting the load into Eq. (13.b), gives:
$\bar{q}(x, y)=\frac{16}{n m \pi^{2}} q_{0} \sin \frac{m \pi A}{2 a} \sin \frac{n \pi B}{2 b}$
(14.d)

And many types of $f(t)$ can be considered as [Khdeir A. A. \& Reddy J. N, 1988]

1. Sine pulse loading
$f(t)=\left\{\begin{array}{ll}\sin \frac{\pi t}{t_{1}} & 0 \leq t \leq t_{1} \\ 0 & t>t_{1}\end{array}\right\}$
2. Step pulse loading
$f(t)=\left\{\begin{array}{lc}1 & 0 \leq t \leq t_{1} \\ 0 & t>t_{1}\end{array}\right\}$
3. Ramp pulse loading
$f(t)=\left\{\begin{array}{cc}\mathrm{t} / t_{1} & 0 \leq t \leq t_{1} \\ 0 & t>t_{1}\end{array}\right\}$

### 2.4 Solution of Dynamic Equilibrium Equations

The equations of equilibrium governing the linear dynamic response of a system can be written as the following formula
$M A \ddot{U}+K S U=R$
(16)

Where:

$$
\begin{gathered}
\mathrm{KS}=\left[\begin{array}{lll}
\hat{c}_{11} & \hat{c}_{12} & \hat{c}_{13} \\
\hat{c}_{12} & \hat{c}_{22} & \hat{c}_{23} \\
\hat{c}_{13} & \hat{c}_{23} & \hat{c}_{33}
\end{array}\right] \quad \mathrm{MA}=\left[\begin{array}{ccc}
\widehat{m}_{11} & 0 & 0 \\
0 & \hat{m}_{22} & 0 \\
0 & 0 & \hat{m}_{33}
\end{array}\right] \quad \ddot{U}=\left\{\begin{array}{c}
\ddot{u}_{m n} \\
\ddot{v}_{m n} \\
\ddot{w}_{m n}
\end{array}\right\} \\
U=\left\{\begin{array}{l}
u_{m n} \\
v_{m n} \\
w_{m n}
\end{array}\right\} \quad R=\left\{\begin{array}{l}
0 \\
0 \\
q
\end{array}\right\}+\left\{\begin{array}{l}
f_{1}^{t} \\
f_{2}^{t} \\
f_{3}^{t}
\end{array}\right\}
\end{gathered}
$$

MA and KS: are the mass and stiffness matrices.
R : is the vector of externally applied loads.
U and $\ddot{\mathrm{U}}$ : are the displacement and acceleration vectors

In the Newmark direct integration method, the first time derivative $\{\dot{\mathrm{U}}\}$ and the solution $\{\mathrm{U}\}$ are approximated at $(\mathrm{n}+1)$ time step (i.e. at time $\mathrm{t}=\mathrm{t}_{\mathrm{n}+1}=(\mathrm{n}+1) \Delta \mathrm{t}$ by the following expression [Rao V. Dukkipati,2010].
$\{\dot{U}\}_{n+1}=\{\dot{U}\}_{n}+\left[(1-\bar{\alpha})\{\ddot{U}\}_{n}+\bar{\alpha}\{\ddot{U}\}_{n+1}\right] \Delta t$
(17)

$$
\begin{equation*}
\{U\}_{n+1}=\{U\}_{n}+\{\dot{U}\}_{n} \Delta t+\left[\left(\frac{1}{2}-\bar{\beta}\right)\{\ddot{U}\}_{n}+\bar{\beta}\{\ddot{U}\}_{n+1}\right](\Delta t)^{2} \tag{18}
\end{equation*}
$$

Where:
$\bar{\alpha}$ And $\bar{\beta}$ : are parameters that control the accuracy and stability of the scheme, and the subscript n indicates that the solution evaluated at $n^{\text {th }}$ time step (i.e. at time, $\mathrm{t}=t_{n}$ ). The choice $\bar{\alpha}=0.5$ and $\bar{\beta}=0.25$ is known to give an unconditionally stable Scheme (average acceleration method), [Rao V. Dukkipati,2010].

## 3. NUMERICAL ANALYSIS

### 3.1 Element Selection and Modeling

An element called shell281 as shown in Fig. 1 is selected which is suitable for analyzing thin to moderately thick shell structures. The element has eight nodes with six degrees of freedom at each node: translations in the $\mathrm{x}, \mathrm{y}$, and z axes, and rotations about the $\mathrm{x}, \mathrm{y}$, and z axes. It may be used for layered applications for modeling composite shells. It is include the effects of transverse shear deformation. The accuracy in modeling composite shells is governed by the first order shear deformation theory. The shell section allows for layered shell definition, options are available for specifying the thickness, material, orientation through the thickness of the layers. But to insert the temperature effect in calculations must be to adding degree of freedom (T). Then, the degrees of freedom change from (6 to7) in each node.

### 3.2 Verification Case Studies

In the present study, Series of preselected cases are modeled to verify the accuracy of the method of analysis. The case study discussed here for dynamic response without temperature change is a comparison of the present work with the numerical solution of [Reddy .J.N, 1982] for a laminated plate Fig. 2,
Close comparison between the two sets of results is evident, for $\mathrm{a} / \mathrm{h}=5$ (maximum central non dimension deflection of present work for CLPT with Newmark direct integration method is= 23.5(error $2.174 \%$ ), for present F.E.M ANSYS maximum central non dimension deflection is $=23.65$ (error $2.8 \%$ ). while for above reference $=23$.
For thermo-mechanical transient response of simply supported laminated plates, the curves of central deflection as a function of time for a ( $0 / 90 / 0$ ) symmetric cross-ply laminated plate subjected to suddenly applied dynamic loading are plotted and compared in Fig. 3 ,with [HuiShen Shen ,2003]. Close comparison between the two sets of results is evident, (maximum central deflection of present work for CLPT with Newmark direct integration method $=2.35 \mathrm{~cm}$
(error2.08\%), and for present F.E.M. ANSYS program is= 2.2375 cm (error6.77\%), while for above reference $=2.4 \mathrm{~cm}$. Fig. 3 .

## 4. EXPERIMENTAL WORK

In the present work, three- purposes were investigated. First, to outline the general steps to design and fabricate the rectangular test models from fiber (E-glass) and polyester resin to form laminate composite materials. Second, the manufactured models are then used to evaluate the mechanical properties ( $E_{1}, E_{2}, G_{12}$ ) with temperature change of unidirectional composite material. Third, evaluate coefficient of thermal expansion (CTE) of the composite plate.

### 4.1 Thermo-Mechanical Analyzer

Thermo-mechanical Analysis (TMA) determines dimensional changes of solids and liquids materials as a function of temperature and/or time under a defined mechanical force.
Irrespective of the selected type of deformation (expansion, compression, penetration, tension or bending), every change of length in the sample is communicated to a highly sensitive inductive displacement transducer (LVDT) via a push rod and transformed into a digital signal. The push rod and corresponding sample holders of fused silica or aluminum oxide can be quickly and easily interchanged to optimize the system to the respective application. Figs. 4 and 5.
The dimension of sample is $(5 * 20 * 4) \mathrm{mm}$. the thermal properties which obtain from this test shown in Table 1.

## 5. RESULTS AND DISCUSSION

The present study focused mainly on the dynamic response behavior of composite laminated plates subjected to mechanical and thermo-mechanical loads of finite duration uniform (step, sine and ramp) and sinusoidal (step, sine and ramp) on the top surface of the plate for three cases of temperature (without temperature effect, $\mathrm{T}=50^{\circ} \mathrm{C}$ and $\mathrm{T}=100^{\circ} \mathrm{C}$ ). The step loading $q(x, y, t)=\bar{q}(x, y), \quad$ ramp loading $\quad q(x, y, t)=\bar{q}(x, y) t / t 1 \quad$ and sinusoid loading $q(x, y, t)=\bar{q}(x, y) \sin \pi t / t 1$. For uniform distributed load $\bar{q}(x, y)=\frac{16}{n m \pi^{2}}$ and for sinusoidal distributed $\bar{q}(x, y)=q_{o}$. The amplitude of force is $q_{o}=100 \mathrm{~N} / \mathrm{mm}^{2}$ and the time of load applied on plate is $t_{1}=0.05 \mathrm{sec}$. The dynamic response of central deflection of composite plate discussed for different parameters such as load condition, lamination angle, temperature change , symmetric or anti symmetric angle ply for simply supported composite plate analytically by CLPT with Newmark direct integration method and numerical result by ANSYS.

## (5-1) Effect of Load Condition

Figs. 6 to 9 represent the variation of central transverse deflection with time (dynamic response) for four layer anti-symmetric and symmetric cross-ply and angle ply simply supported laminated plates under sinusoidal $\quad\left(P(x, y)=q o \sin \left(\pi_{x} / a\right) \sin \left(\pi_{y} / b\right)\right)$ and uniform $(P(x, y)=q o)$ variation loading, (step $q(x, y, t)=P(x, y)$, ramp loading $q(x, y, t)=$ $P(x, y) t / t 1$ and sinusoid loading $\left.\mathrm{q}(\mathrm{x}, \mathrm{y}, \mathrm{t})=\mathrm{P}(\mathrm{x}, \mathrm{y}) \sin \pi \mathrm{t} / \mathrm{t}_{1}\right)$ for $q o=100 \mathrm{~N} / m^{2}, \mathrm{t}_{1}=0.05 \mathrm{sec}$ ) without any temperature change solved analytically by CLPT with Newmark direct integration method and (F.E.M) by ANSYS program. The deflection due to step loading higher in magnitude than the other loads with percentage reach to $91.96 \%, 97.4 \%$ from sine and ramp load, respectively, because the step load subjected suddenly with constant value with the time.

Very good verification between CLPT with Newmark and FEM by ANSYS maximum error is $12.9 \%$. Maximum response for step load always occurs in the time of applying load (i.e. in the time less than $t_{1}$ after that the response became in negative sign and positive sign alternatively. for ramp load, the response increasing linearly with time until it reached to $t_{1}$ at this point the maximum response occurs, then the response became in negative sign and positive sign alternatively. For sine load the response behavior have the sine shape and the maximum response at $t_{1} / 2$.

## (5-2) Effect of Temperature Change with Varies Load Condition

Figs. 10 to 13 show the numerical result by ANSYS for dynamic response of central deflection of symmetric and anti-symmetric angle ply simply supported composite plate step uniform and step sinusoidal load and different condition of temperature effect i.e. $\left(\mathrm{T}=25^{\circ} \mathrm{C}, \mathrm{T}=50^{\circ} \mathrm{C}\right.$, $\mathrm{T}=100^{\circ} \mathrm{C}$ ). The deflection increases with percentage reaches to ( $58.47 \%$ ) when temperature became $50^{\circ} \mathrm{C}$ and when the temperature reach to $100^{\circ} \mathrm{C}$ the response increases with higher percentage reaches to ( $200 \%$ ) with respect to response without change in temperature for laminated plates for step uniform dynamic load.

The reason behind that is there where two loads (mechanical and thermal) each load causes the deflections (thermal and mechanical deflections) summation is the deflection of plate under thermo-mechanical loading. When the temperature increases the deflection increased with high percentage. The uniform load is higher than sinusoidal load for all load condition.

## (5-3) Effect of Lamination Angle

Fig. 14 shows the effect of angle $(\theta)$ on central deflection for four layer symmetric angle-ply laminated plates, simply supported, subjected to sine uniform loading with applied temperature equal to $50^{\circ} \mathrm{C}$, solved analytically by Newmark and numerical by ANSYS .From the results, the central deflection of laminated plate decreases with increasing the angle ( $\theta$ ) from 10 to 40 with percentage reach to $24.8 \%$. Then increase the central deflection when $\theta$ increase from 40 to 70 with percentage reaches to $9.2 \%$. The maximum deflection with time for each case is when lamination angle is 10 .

## 6. CONCLUSION

This study considers the vibration analysis of symmetric and anti-symmetric angle-ply composite laminate plate. From the present study, the following conclusions can be made:
1-The Young and shear modulus decrease when temperature increases with high percentages reach to $96.3 \%$ when temperature changes from $\left(20{ }^{\circ} \mathrm{C}\right.$ to $\left.100^{\circ} \mathrm{C}\right)$ for longitudinal young modulus, for transverse young modulus is $96.53 \%$ and for shear modulus is $91.1 \%$. The longitudinal and transverse coefficient of thermal expansion also decrease when temperature increase with percentage $80 \%$ and $73.7 \%$ respectively for the same temperature.
2- The response due to step loading higher in magnitude than the other loads with percentage reach to $91.96 \%, 97.4 \%$ from sine and ramp load, respectively.

3- The response increase with maximum percentage reaches to (58.47\%) when temperature became $50^{\circ} \mathrm{C}$ and when the temperature reach to $100^{\circ} \mathrm{C}$ the response increase with higher percentage reaches to $(200 \%)$ with respect to response without change in temperature 4- It was seen that the different fiber orientation angles affected on dynamic response. The central deflection of laminated plate decreases with increasing the angle $(\theta)$ from 10 to 40 with percentage reach to $24.8 \%$. Then increase the central deflection when $\theta$ increase from 40 to 70 with percentage reaches to $9.2 \%$. Thus, the maximum deflection with time is when lamination angle is 10 for four layer symmetric angle-ply laminated plates, simply supported, subjected to sine uniform loading with applied temperature equal to $50^{\circ} \mathrm{C}$.

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## NOMENCLATURE

| Symbol | Description | Unit |
| :---: | :---: | :---: |
| a, b | Dimension of plate in x and y coordinate | m |
| $\mathrm{A}_{\mathrm{ij}}, \mathrm{B}_{\mathrm{ij}}, \mathrm{D}_{\mathrm{ij}}$ | Extensional stiffness, the coupling stiffness, and the bending stiffness | - |
| $\mathrm{E}_{1}, \mathrm{E}_{2}, \mathrm{E}_{3}$ | Elastic modulus of composite material | GPa |
| $\mathrm{G}_{12}, \mathrm{G}_{23}, \mathrm{G}_{13}$ | Shear modulus of composite material | GPa |
| h | Thickness | m |
| $I_{0}, I_{1}, I_{2}$ | Mass moment of inertia | kg.m ${ }^{2}$ |
| [MA] | Mass matrix | kg |
| $M_{x x}, M_{y y}, M_{x y}$ | Moment resultant per unit length | N.m/m |
| N | Total number of plate layers | - |
| Nxx, Nyy , Nxy | The resultant of in-plane force per unit length | N/m |
| $N_{x x}^{t}, N_{y y}^{t}, N_{x y}^{t}$ | The resultant of in-plane force per unit length with thermal effect | N/m |
| $\widehat{N} x x, \widehat{N} y y, \widehat{N} x y$ | Applied edge force | N/m |
| $\mathrm{q}(\mathrm{x}, \mathrm{y}, \mathrm{t})$ | Dynamic force subjected on a system | $\mathrm{N} / \mathrm{m}^{2}$ |
| $\bar{Q}_{i j}^{(k)}$ | Transformed lamina stiffness | N/m |
| R | Vector of externally applied loads | N |
| t | Time | min or s |
| $\Delta \mathrm{t}$ | Time Interval | min or s |
| $t_{1}$ | The end time of load | sec |
| T | Temperature | $\mathrm{C}^{0}$ |
| $\Delta \mathrm{T}$ | Temperature increment | $\mathrm{C}^{0}$ |


| $T_{\text {ref }}$ | Reference temperature | $\mathrm{C}^{0}$ |
| :---: | :---: | :---: |
| $U, \dot{U}, \ddot{U}$ | Displacement , velocity and acceleration vectors | $\mathrm{m}, \mathrm{m} / \mathrm{s}, \mathrm{m} / \mathrm{s}^{2}$ |
| $u_{o}, v_{o}, w_{o}$ | Displacement components along (x,y,z) directions <br> respectively | m |
| $U_{m n}, V_{m n}, W_{m n}$ | Amplitude of $\left(u_{o}, v_{o}, w_{o}\right)$ respectively | - |
| $\mathrm{x}, \mathrm{y}, \mathrm{z}$ | Cartesian coordinate system | m |
| z | Distance from neutral axis | m |
| $\theta$ | Fiber orientation angle | Degree |
| $\alpha_{1}, \alpha_{2}$ | Coefficient of thermal |  |
| expansion of composite material | $\left(1 / \mathrm{K}^{-1}\right)$ |  |
| $\rho$ | Density | $(\mathrm{kg} / \mathrm{m} 3)$ |
| $\epsilon_{x x}, \epsilon_{y y}, \epsilon_{x y}$ | Strain components | $\mathrm{m} / \mathrm{m}$ |
| $\sigma_{x x}, \sigma_{y y}, \sigma_{x y}$ | Stress components | GPa |



Figure 1. Shell281 geometry [ANSYS 13 Program].


Figure 2. Comparison of the present solution with the numerical solution of [Reddy .J.N, 1982] of two-layer cross-ply ( $0 / 90$ ) square plate under suddenly applied sinusoidal loading $(\mathrm{a} / \mathrm{h}=5$ ).



Figure 3. Comparison of present study with [Hui-Shen Shen et al 2003] for laminated square plate under thermal loading condition at $\left(\Delta T=200 \mathrm{C}^{0}\right)$.


Figure 4. Operating principle of TMA.


Figure 5. TMA PT1000 device.

Table1. Experimental value of mechanical and thermal properties of fiber -polyester composite plate for fiber volume fraction $=0.3$ changed with temperature.

| $\mathbf{T}$ <br> $\mathbf{C}^{\circ}$ | $\boldsymbol{E}_{\mathbf{1}}$ <br> $\mathbf{M p a}$ | $\boldsymbol{E}_{\mathbf{2}}$ <br> $\mathbf{M p a}$ | $\boldsymbol{G}_{\mathbf{1 2}}=\boldsymbol{G}_{\mathbf{1 3}}=\boldsymbol{G}_{\mathbf{2 3}}$ <br> $\mathbf{M p a}$ | $\boldsymbol{\alpha}_{\mathbf{1}}$ <br> $\mathbf{E - 6 / K}$ | $\boldsymbol{\alpha}_{\mathbf{2}}$ <br> $\mathbf{E - 6 / K}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{2 0}$ | 24627.0 | 5588.04 | 1551.77 | 14.57 | 47.81 |
| $\mathbf{3 0}$ | 23343.30 | 4123.11 | 1618.65 | 9.03 | 31.7 |
| $\mathbf{4 0}$ | 21775.40 | 1550.22 | 2505.75 | 7.20 | 29.36 |
| $\mathbf{5 0}$ | 15219.80 | 1515.37 | 623.423 | 4.79 | 25.79 |
| $\mathbf{6 0}$ | 6475.41 | 566.8 | 114.2336 | 3.20 | 21.38 |
| $\mathbf{7 0}$ | 2990.82 | 458.59 | 113.535 | 3.18 | 15.60 |
| $\mathbf{8 0}$ | 2555.71 | 289.27 | 130.48 | 3.22 | 15.59 |
| $\mathbf{9 0}$ | 1471.90 | 210.49 | 158.83 | 3.08 | 15.19 |
| $\mathbf{1 0 0}$ | 903.90 | 193.84 | 138.48 | 2.91 | 12.58 |
| $\mathbf{1 1 0}$ | 741.31 | 191.75 | 131.74 | 2.75 | 11.57 |
| $\mathbf{1 2 0}$ | 674.40 | 187.53 | 125.23 | 2.57 | 10.47 |
| $\mathbf{1 3 0}$ | 644.70 | 186.51 | 122.56 | 2.48 | 9.34 |
| $\mathbf{1 4 0}$ | 629.01 | 185.19 | 117.59 | 2.45 | 7.67 |
| $\mathbf{1 5 0}$ | 612.02 | 182.66 | 107.81 | 2.44 | 5.28 |
| $\mathbf{1 6 0}$ | 597.61 | 181.88 | 100.28 | 2.44 | 4.19 |
| $\mathbf{1 7 0}$ | 592.41 | 173.50 | 95.414 | 2.44 | 3.90 |
| $\mathbf{1 8 0}$ | 592.22 | 164.55 | 95.04 | 2.45 | 3.88 |
| $\mathbf{1 9 0}$ | 591.02 | 163.91 | 85.71 | 2.46 | 3.66 |
| $\mathbf{2 0 0}$ | 590.57 | 153.3 | 83.64 | 2.47 | 3 |



Figure 6. Central deflection of four layers symmetric angle-ply (45/-45/...) laminated plates for variant sinusoidal dynamic load without temperature change.


Figure 7. Central deflection of four layers symmetric angle-ply (45/-45/...) laminated plates for variant uniform dynamic load without temperature change.


Figure 8. Central deflection of four layers anti-symmetric angle-ply (45/-45/...) laminated plates for variant sinusoidal dynamic load without temperature change.


Figure 9. Central deflection of four layers anti-symmetric angle-ply (45/-45/...) laminated plates for variant uniform dynamic load without temperature change.


Figure 10. Central deflection of four layers symmetric angle-ply (45/-45/...) laminated plates for step uniform dynamic load with temperature change.


Figure 11. Central deflection of four layers symmetric angle-ply (45/-45/...) laminated plates for step sinusoidal dynamic load with temperature change.


Figure 12. Central deflection of four layers anti- symmetric angle-ply (45/-45/...) laminated plates for step uniform dynamic load with temperature change.


Figure 13. Central deflection of four layers anti- symmetric angle -ply (45/-45/...) laminated plates for step sinusoidal dynamic load with temperature change.

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Figure 14. Effect of lamination angle on central deflection of four layers $(\theta /-\theta /-\theta / \theta)$ laminated plates for sine uniform dynamic load with temperature equal to $50^{\circ} \mathrm{C}$.

